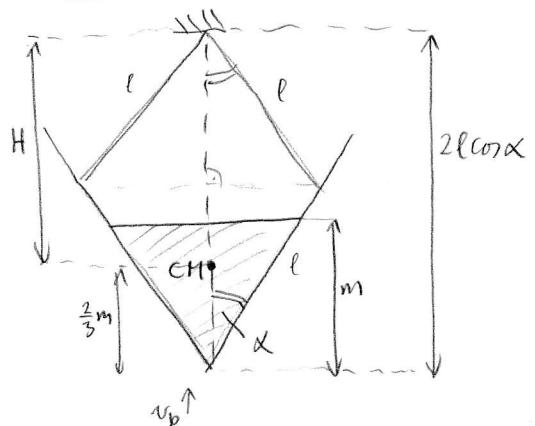


- The equilibrium-point of the system can be found from the potential energy of the system which is minimal (for a given amount of water).

figure 1:



- the total volume of water: $V = L m^2 \operatorname{tg} \alpha \Rightarrow m = \sqrt{\frac{V}{L \operatorname{tg} \alpha}} = \sqrt{\frac{S}{\operatorname{tg} \alpha}}$
where $\frac{V}{L} \equiv S$ is constant

- the distance of the centre of mass from the ceiling:

$$H = 2l \cos \alpha - \frac{2}{3} m = 2l \cos \alpha - \frac{2}{3} \sqrt{\frac{S}{\operatorname{tg} \alpha}}$$

- the gravitational potential energy is minimal when $\left. \frac{dH}{d\alpha} \right|_{\alpha=\alpha_0} = 0$

$$(1) -2l \sin \alpha_0 + \frac{1}{3} \sqrt{S \operatorname{tg} \alpha_0} \left(1 + \frac{1}{\operatorname{tg}^2 \alpha_0} \right) = 0$$

- the water almost fills up the entire vessel, so:

$$(2) \frac{V}{L} = S = l^2 \sin \alpha_0 \cos \alpha_0$$

- (1) and (2): $-2l \sin \alpha_0 + \frac{1}{3} \left[l^2 \sin \alpha_0 \cos \alpha_0 \cdot \frac{\sin \alpha_0}{\cos \alpha_0} \right] \left(1 + \frac{1}{\operatorname{tg}^2 \alpha_0} \right) = 0$

$$-2l \sin \alpha_0 + \frac{1}{3} l \sin \alpha_0 \left(1 + \frac{1}{\operatorname{tg}^2 \alpha_0} \right) = 0 \quad ! \sin \alpha_0 \neq 0$$

$$-2 + \frac{1}{3} \left(1 + \frac{1}{\operatorname{tg}^2 \alpha_0} \right) = 0$$

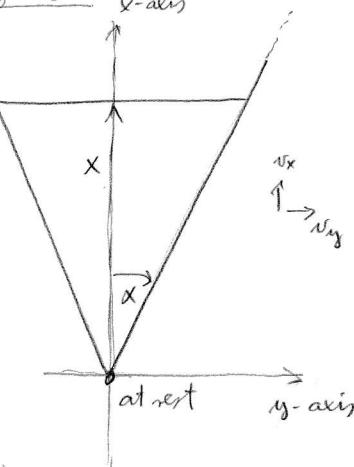
$$\operatorname{tg} \alpha_0 = \frac{1}{\sqrt{5}}$$

$$! 0 < \alpha_0 < \frac{\pi}{2}$$

The angle between the plates: $2\alpha_0 = 2 \cdot \arctg \left(\frac{1}{\sqrt{5}} \right) \approx 0,841 \text{ rad} \approx 48,19^\circ$

- Now let us relate the speed of the water particles to the angular velocity of the vessel ($\dot{\alpha}$). Let the bottom point of the vessel be at rest. According to the first hint if we draw a set of equispaced horizontal lines onto the water, they preserve these properties. This indicates that the vertical velocity of water particles are the same in a horizontal plane. The volume of water under such a line is constant.

figure 2:



- $\frac{V(x)}{L} = x^2 \operatorname{tg} \alpha = S(x)$ is constant

- derivating with respect to time (and taking values around the c. p.)

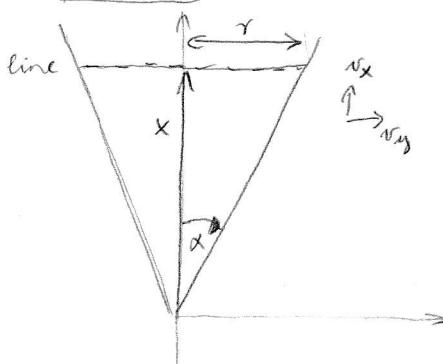
$$2x \cdot \frac{dx}{dt} \operatorname{tg} \alpha_0 + x^2 \cdot \frac{1}{\cos^2 \alpha_0} \cdot \frac{d\alpha}{dt} = 0$$

$$v_x(x) = \frac{dx}{dt} = -\frac{1}{2} \frac{1}{\operatorname{tg} \alpha_0 \cos^2 \alpha_0} \cdot x \cdot \dot{\alpha} = -\frac{3}{\sqrt{5}} x \cdot \dot{\alpha}$$

$$\left(\operatorname{tg} \alpha_0 = \frac{1}{\sqrt{5}} \Rightarrow \sin \alpha_0 = \frac{1}{\sqrt{6}} ; \cos \alpha_0 = \frac{\sqrt{5}}{\sqrt{6}} \right)$$

- As the oscillation is symmetric to mirror-symmetric, the horizontal velocity of water particles at $y=0$ is zero. Water particles near the plates follow its motion horizontally, so horizontal velocity at these point also can be determined:

figure 3:



$$\frac{V}{L} = x^2 \operatorname{tg}\alpha = \frac{r^2}{\operatorname{tg}\alpha} = S(x) = \text{const.} \quad (r = x \cdot \operatorname{tg}\alpha)$$

taking derivative of both sides with respect to time:

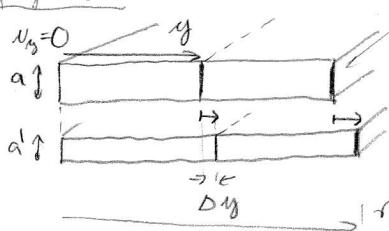
$$\frac{2r \frac{dr}{dt} \operatorname{tg}\alpha_0 - r^2 \cdot \frac{1}{\cos^2 x_0} \frac{dx}{dt}}{\operatorname{tg}^2 \alpha} = 0 \quad (\operatorname{tg}\alpha \neq 0)$$

$$\frac{dx}{dt} = v_y(x, x \operatorname{tg}\alpha) = +\frac{1}{2} \frac{1}{\operatorname{tg}\alpha_0 \cos^2 x_0} \cdot r \dot{\alpha} = +\frac{3}{\sqrt{5}} (x \operatorname{tg}\alpha_0) \dot{\alpha}$$

(holds true for the $y \geq 0$ region)

- Considering the water between two adjacent horizontal line one can easily derive that the horizontal velocity must increase linearly with the distance from x-axis (the two boundary plates remain parallel and there is no water flux through the plates)

figure 4:



$$\alpha y L = \text{const} \Rightarrow \frac{da}{dt} y + a \frac{dy}{dt} = 0 \Rightarrow \frac{dy}{dt} = -\frac{1}{a} \frac{da}{dt} \alpha L y$$

$$v_y(x, 0) = 0; v_y(x, y = x \operatorname{tg}\alpha) = +\frac{3}{\sqrt{5}} (x \operatorname{tg}\alpha_0) \dot{\alpha} \Rightarrow v_y(x, y) = +\frac{3}{\sqrt{5}} y \dot{\alpha} \quad (\text{holds true for the entire region})$$

- The bottom of the vessel moves vertically with respect to the ceiling, its speed is:

$$v_b = -\frac{d}{dt}(2L \cos \alpha) = +2L \sin \alpha_0 \cdot \dot{\alpha} = \frac{2}{\sqrt{6}} L \dot{\alpha}$$

- So the velocity of a water particle with respect to the ceiling:

$$v_x(x) = v_b + v_x(x) = \frac{2}{\sqrt{6}} L \dot{\alpha} - \frac{3}{\sqrt{5}} x \dot{\alpha}; v_y(x, y) = +\frac{3}{\sqrt{5}} y \dot{\alpha}$$

- The total kinetic energy of the water:

$$dK = \frac{1}{2} (L dx dy \beta) (v_x^2(x, y) + v_y^2(x, y)) = \frac{1}{2} (L dx dy \beta) \left[\left(\frac{2}{\sqrt{6}} L \dot{\alpha} - \frac{3}{\sqrt{5}} x \dot{\alpha} \right)^2 \dot{x}^2 + \left(\frac{3}{\sqrt{5}} y \dot{\alpha} \right)^2 \dot{y}^2 \right]$$

$$K = 2 \int_0^{\frac{\pi}{2} \cos \alpha_0} dx \int_0^{x \operatorname{tg} \alpha_0} dy \left\{ \frac{1}{2} L \beta \dot{\alpha}^2 \left[\left(\frac{2}{\sqrt{6}} L \dot{\alpha} - \frac{3}{\sqrt{5}} x \dot{\alpha} \right)^2 + \left(\frac{3}{\sqrt{5}} y \dot{\alpha} \right)^2 \right] \right\} = L \beta \dot{\alpha}^2 \int_0^{\frac{\pi}{2} L} dx \int_0^{\frac{1}{\sqrt{5}} x} dy \left\{ \left(\frac{2}{\sqrt{6}} L \dot{\alpha} - \frac{3}{\sqrt{5}} x \dot{\alpha} \right)^2 + \left(\frac{3}{\sqrt{5}} y \dot{\alpha} \right)^2 \right\} = L \beta \dot{\alpha}^2 \int_0^{\frac{\pi}{2} L} \left\{ \left(\frac{2}{\sqrt{6}} L \dot{\alpha} - \frac{3}{\sqrt{5}} x \dot{\alpha} \right)^2 \frac{x}{\sqrt{5}} + \frac{3}{25 \sqrt{5}} x^3 \right\} dx = L \beta \dot{\alpha}^2 L^4 \frac{\sqrt{5}}{90}$$

- Let us introduce $\xi := x - x_0$
- The total kinetic energy of the water: $K = L \dot{\xi}^4 \cdot \frac{\sqrt{5}}{30} \cdot \dot{\xi}^2$
- Now let us express the gravitational potential energy around the equilibrium point as a function of ξ ($\xi \ll l$) up to second order (we choose the potential energy zero at the e. point):
 - distance of the centre of mass from the ceiling: $H = 2l \cos \alpha - \frac{2}{3} \sqrt{\frac{S}{L}} \alpha$ (S is constant)
 - $\frac{dH}{dx} \Big|_{x=x_0} = 0$, $\frac{d^2H}{dx^2} \Big|_{x=x_0} = -\frac{2}{5} \sqrt{30} \cdot l$

$$\Rightarrow \Pi(\xi) = -H(\xi)Mg - (-H(\xi=0)Mg) = -\left(H(0) + \frac{dH}{dx} \Big|_{x=x_0} \cdot \xi + \frac{1}{2} \frac{d^2H}{dx^2} \Big|_{x=x_0} \xi^2\right)Mg + H(0)Mg = +\frac{2}{5} \sqrt{30} l \xi^2 Mg$$

where $M = (SL)S = l^2 \sin \alpha \cos \alpha L S = \frac{\sqrt{5}}{6} l^2 L S$

- The total energy of the system (kinetic + potential):

$$E = K + \Pi = \frac{\sqrt{5}}{30} L \dot{\xi}^4 \cdot \dot{\xi}^2 + \frac{\sqrt{6}}{3} L \dot{\xi}^3 \cdot g \cdot \dot{\xi}$$

- If there is no friction in the water and between the water and the plates, E is conserved. Taking time derivatives of both sides:

$$0 = \frac{\sqrt{5}}{30} L \dot{\xi}^4 \cdot 2\ddot{\xi}\dot{\xi} + \frac{\sqrt{6}}{3} L \dot{\xi}^3 \cdot 2g\dot{\xi}$$

$$0 = \frac{2\sqrt{5}}{30} l \cdot \ddot{\xi} + \frac{2\sqrt{6}}{3} g \cdot \dot{\xi}$$

$$\ddot{\xi} = -\frac{g}{l} \cdot 6\sqrt{30} \cdot \dot{\xi}$$

angular frequency of the oscillation: $\omega = \sqrt{\frac{g}{l}} \cdot \sqrt{6\sqrt{30}} \approx 5,733 \sqrt{\frac{g}{l}}$