

WoPhO Selection Round Problem 2

Two-ball body

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Part A

$$1 \quad x_0 = \frac{m_1 \cdot 0 + m_2 \cdot (R_1 + R_2)}{m_1 + m_2} = \frac{\frac{4\pi}{3}\rho R_2^3 (R_1 + R_2)}{\frac{4\pi}{3}\rho(R_1^3 + R_2^3)} = \frac{R_2^3}{R_1^3 - R_1 R_2 + R_2^2} = 5.33 \text{ cm.}$$

2 This axis is a CM axis of both balls, therefore

$$I_{\parallel} = \frac{2}{5}m_1 R_1^2 + \frac{2}{5}m_2 R_2^2 = \frac{8\pi}{15}\rho(R_1^5 + R_2^5) = 2.654 \cdot 10^{-4} \text{ kg m}^2.$$

3 The parallel CM axes of the spheres are respectively x_0 and $x'_0 = R_1 + R_2 - x_0 = \frac{R_1^3}{R_1^2 - R_1 R_2 + R_2^2}$ away from the examined axis, thus according to Steiner's theorem

$$I_{\perp} = \frac{2}{5}m_1 R_1^2 + m_1 x_0^2 + \frac{2}{5}m_2 R_2^2 + m_2 x_0'^2 = \frac{8\pi}{15}\rho(R_1^5 + R_2^5) + \frac{4\pi}{3}\rho \frac{R_1^3 R_2^3 (R_1 + R_2)}{R_1^2 - R_1 R_2 + R_2^2} = 4.263 \cdot 10^{-4} \text{ kg m}^2.$$

Part B

1 The body will constitute pure rolling, therefore the contact point with the ground will always be a momentary rotation center of the motion. The kinetic energy of the body is therefore $\frac{1}{2}I\omega^2$ where I is the moment of inertia with respect to that point. Now we're going to calculate this moment of inertia at the moment of the collision. The cosine of α in the figure is clearly $\cos \alpha = -\sin(\alpha - 90^\circ) = -\frac{R_2 - R_1}{R_2 + R_1} = \frac{R_1 - R_2}{R_1 + R_2}$. From the law of cosines we get for the distance of the CM and the rotation center $d^2 = R_1^2 + x_0^2 - 2R_1 x_0 \cos \alpha$; by Steiner's theorem

$$I = I_{\perp} + (m_1 + m_2)d^2 = \frac{4\pi}{15}\rho(7R_1^5 + 7R_2^5 + 20R_1 R_2^4).$$

The kinetic energy is on the other hand the change of gravitational potential energy of the system: that of the lower ball doesn't vary as its CM moves only horizontally, but the potential energy of the upper ball changes from (with choosing $V = 0$ at ground level) $m_2 g(2R_1 + R_2)$ to $m_2 g R_2$. By the conservation of energy:

$$m_2 g(2R_1 + R_2) = m_2 g R_2 + \frac{1}{2}I\omega^2$$

$$\omega = \sqrt{\frac{4m_2 g R_1}{I}} = \sqrt{\frac{20R_2^3 R_1 g}{7R_1^5 + 7R_2^5 + 20R_1 R_2^4}}.$$

The velocities of the balls' centers can be found from the fact that the ground contact point is the momentary center of rotation, therefore the velocity of each point of the body is $r\omega$ where r is the distance of this point and the ground contact point and the velocity vector is perpendicular to the line joining them (this is shown in the sketch). The distance of the center of ball 1 to the ground contact point is R_1 , thus its velocity is

$$v_1 = R_1 \omega = \sqrt{\frac{20R_2^3 R_1^3 g}{7R_1^5 + 7R_2^5 + 20R_1 R_2^4}} = 0.2386 \frac{\text{m}}{\text{s}};$$

that for ball 2 is $r_2 = \sqrt{R_1^2 + (R_1 + R_2)^2 - 2R_1(R_1 + R_2) \cos \alpha} = \sqrt{4R_1 R_2 + R_2^2}$, thus

$$v_2 = r_2 \omega = \sqrt{\frac{20R_2^4 R_1 (4R_1 + R_2) g}{7R_1^5 + 7R_2^5 + 20R_1 R_2^4}} = 0.8264 \frac{\text{m}}{\text{s}}.$$

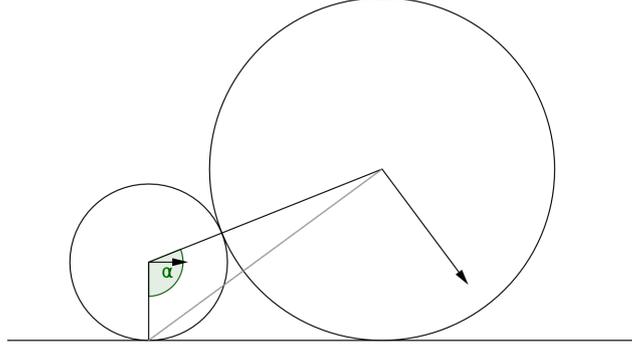


Figure 1. The setup of the system immediately before the collision in Part B.1.

2 As there is no friction in the system, there is no horizontal force at all, therefore the horizontal component of the system's momentum is conserved and since it was 0 initially, it will be 0 in the final setup as well. Consequently, the horizontal component of the CM's velocity is 0, thus it will move vertically. The center of the lower ball must have a purely horizontal velocity, as the contact point with the ground can only move so and its relative velocity to the mentioned center is horizontal too (it is a peripheral speed related to a vertical position vector).

The vector from the CM to the center of the lower ball in the final state (Figure 2) is $\mathbf{r}_1 = x_0(-\sin \alpha, \cos \alpha)$, thus the vector of the relative velocity is $\mathbf{v}_{\text{rel},1} = \boldsymbol{\omega} \times \mathbf{r}_1 = x_0\omega(\cos \alpha, \sin \alpha)$. The only way to decompose this into the difference of a horizontal and a vertical velocity vector (as required by the above conditions) gives the velocity of the CM as $v_{\text{CM}} = -x_0\omega \sin \alpha$ (and that of the center of the lower ball as $v_1 = x_0\omega \cos \alpha$).

The kinetic energy of the system right before the collision is given as the sum of the translational kinetic energy due to v_{CM} and the rotational energy: this must be equal to the reduction of potential energy, which is $2m_2gR_1$ (see above):

$$\frac{1}{2}(m_1 + m_2)v_{\text{CM}}^2 + \frac{1}{2}I_{\perp}\omega^2 = 2m_2gR_1$$

$$\omega^2 (x_0^2\omega^2 \sin^2 \alpha(m_1 + m_2) + I_{\perp}) = 4m_2gR_1$$

$$\omega = \sqrt{\frac{4m_2gR_1}{x_0^2\omega^2 \sin^2 \alpha(m_1 + m_2) + I_{\perp}}} = \sqrt{\frac{20R_2^3R_1(R_1^3 + R_2^3)g}{2R_1^8 + 7R_1^5R_2^3 + 10R_1^4R_2^4 + 7R_1^3R_2^5 + 20R_1R_2^7 + 2R_2^8}}$$

(We have used that $\sin^2 \alpha = 1 - \cos^2 \alpha = \frac{4R_1R_2}{(R_1+R_2)^2}$.)

The velocity of the center of the lower ball is horizontal and its magnitude is

$$v_1 = x_0\omega \cos \alpha = \sqrt{\frac{20R_2^9R_1(R_1 - R_2)^2g}{(2R_1^8 + 7R_1^5R_2^3 + 10R_1^4R_2^4 + 7R_1^3R_2^5 + 20R_1R_2^7 + 2R_2^8)(R_1^3 + R_2^3)}} = 0.2519 \frac{\text{m}}{\text{s}}$$

The distance vector from the center of ball 1 to that of ball 2 is $\mathbf{d} = (R_1 + R_2)(\sin \alpha, -\cos \alpha)$ so the relative velocity of the center of ball 2 to that of ball 1 is $\mathbf{v}_{\text{rel}} = \boldsymbol{\omega} \times \mathbf{d} = (R_1 + R_2)\omega(-\cos \alpha, -\sin \alpha)$; since the velocity vector of the center of ball 1 is $\mathbf{v}_1 = (x_0\omega \cos \alpha, 0)$, that of ball 2 is

$$\mathbf{v}_2 = \omega(-x_0' \cos \alpha, -(R_1 + R_2) \sin \alpha).$$

The angle β between this vector and the positive x -axis is given by

$$\tan \beta = \frac{v_{2y}}{v_{2x}} = \frac{R_1 + R_2}{x_0'} \tan \alpha = \frac{R_1^3 + R_2^3}{R_1^3} \cdot \frac{2\sqrt{R_1R_2}}{R_1 - R_2} \rightarrow \beta = \arctan \left(\frac{2\sqrt{R_1R_2}}{R_1 - R_2} \frac{R_1^3 + R_2^3}{R_1^3} \right) = -87.75^\circ,$$

while its magnitude is

$$\begin{aligned} v_2 &= |\mathbf{v}_2| = \omega \sqrt{x_0'^2 \cos^2 \alpha + (R_1 + R_2)^2 \sin^2 \alpha} \\ &= \sqrt{\frac{20R_2^3R_1^2g(R_1^7 + 2R_1^6R_2 + R_1^5R_2^2 + 8R_1^3R_2^4 + 4R_2^7)}{(2R_1^8 + 7R_1^5R_2^3 + 10R_1^4R_2^4 + 7R_1^3R_2^5 + 20R_1R_2^7 + 2R_2^8)(R_1^3 + R_2^3)}} = 0.8022 \frac{\text{m}}{\text{s}}. \end{aligned}$$

The specialities of the velocities are indicated in Figure 2. (Note that the momentary center of rotation will be the intersection of the vertical through the center of ball 1 and the horizontal through the CM: this follows from the directions of the velocities of the center of ball 1 and the CM.)

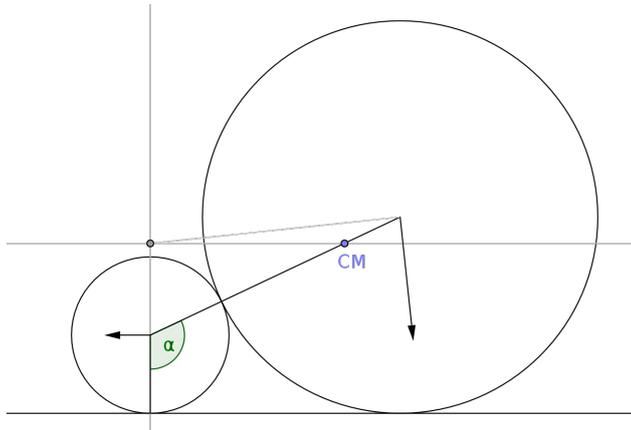


Figure 2. The setup of the system immediately before the collision in Part B.2.

Part C

Both contact points with the ground will move without slipping. Consequently, if the body rotates about its symmetry axis by some angle, these points will move by distances the ratio of which is R_2/R_1 : this means that the orbits of these points are two concentric circles with radii of this ratio. The common center of these circles will be a permanent center of the rotation, therefore we will calculate with respect to axes passing through this point. Now we determine the distance D of this center to the CM. As the center O must be the intersection point of the symmetry axis O_1O_2 and the plane the body is placed on, we can write down the following equation due to the similar right triangles OO_1T_1 and OO_2T_2 :

$$\frac{D + x'_0}{R_2} = \frac{D - x_0}{R_1} \rightarrow D = \frac{R_1^4 + R_2^4}{(R_1^2 - R_1R_2 + R_2^2)(R_2 - R_1)}.$$

However, the line joining the touching points of the body and the plane is the momentary rotational axis of the body due to the slippless motion. Therefore, the angular velocity vector ω must be in this line. The angle γ in the figure can simply be expressed by $\sin \gamma = \frac{R_2 - R_1}{R_2 + R_1} \rightarrow \cos \gamma = \frac{2\sqrt{R_1R_2}}{R_1 + R_2}$, so the components of this vector in the principal coordinate system are $\omega_1 = \omega \sin \gamma$ and $\omega_2 = \omega \cos \gamma$. The moments of inertia with respect to axes passing through this point are $I_\perp + (m_1 + m_2)D^2$ and I_\parallel , respectively.

Now we're going to write down the total energy of the system. The kinetic energy can be written solely as rotational energy with respect to point O : using the above expressions of the moments of inertia and angular velocities we find

$$K = \frac{1}{2}(I_\perp + (m_1 + m_2)D^2)\omega^2 \sin^2 \gamma + \frac{1}{2}I_\parallel\omega^2 \cos^2 \gamma = \frac{14\pi}{15}\rho\omega^2(R_1^5 + R_2^5).$$

The CM of the system is moving on a circle in a plane somewhat above the inclined plane and parallel to it. The radius of this circle is $R = D \cos \gamma = \frac{2(R_1^4 + R_2^4)\sqrt{R_1R_2}}{(R_1^2 - R_1R_2 + R_2^2)(R_2 - R_1)}$. Let θ be the angle between the line joining the touching points and the horizontal edge of the plane. If we choose the zero level of the gravitational potential energy at the state $\theta = \pi/2$, we can simply express the potential energy as

$$V = (m_1 + m_2)g\Delta h = -(m_1 + m_2)gR \cos \theta \sin \alpha = \frac{8\pi}{3}\rho g \frac{R_1^4 + R_2^4}{R_1 - R_2} \sqrt{R_1R_2} \cos \theta \sin \alpha.$$

As there is no energy loss due to friction, the total energy $K + V$ is constant.

1 The energy initially is 0 as $\omega = 0$ and $\theta = \pi/2$: due to the conservation of energy, this will hold all along the motion. Therefore, the less is the potential energy, the more is the kinetic energy and so the velocities. From the expression of V we can see that the minimal potential energy at $\theta = 0$ is $\frac{8\pi}{3}\rho g \frac{R_1^4 + R_2^4}{R_1 - R_2} \sqrt{R_1R_2} \sin \alpha$, the opposite of which is the maximal kinetic energy:

$$\frac{8\pi}{3}\rho g \frac{R_1^4 + R_2^4}{R_2 - R_1} \sqrt{R_1R_2} \sin \alpha = \frac{14\pi}{15}\rho\omega_{\max}^2(R_1^5 + R_2^5)$$

$$\omega_{\max} = \sqrt{\frac{20}{7} g \frac{(R_1^4 + R_2^4) \sqrt{R_1 R_2}}{(R_2 - R_1)(R_1^5 + R_2^5)}} \sin \alpha;$$

as the distances of the centers of the balls to the rotational axis are R_1 and R_2 , the maximal velocities of these centers are

$$v_{1,\max} = R_1 \sqrt{\frac{20}{7} g \frac{(R_1^4 + R_2^4) \sqrt{R_1 R_2}}{(R_2 - R_1)(R_1^5 + R_2^5)}} \sin \alpha = 0.4517 \frac{\text{m}}{\text{s}};$$

$$v_{2,\max} = R_2 \sqrt{\frac{20}{7} g \frac{(R_1^4 + R_2^4) \sqrt{R_1 R_2}}{(R_2 - R_1)(R_1^5 + R_2^5)}} \sin \alpha = 0.9033 \frac{\text{m}}{\text{s}}.$$

2 The magnitude of ω is maximal in this point, therefore it doesn't change in the first order. However, its direction is varying over time, and this causes an angular acceleration like in the case of centripetal acceleration. We are now going to calculate the angular velocity $\Omega = \dot{\theta}$ of the rotation of the axis $T_1 T_2$: as its vector is perpendicular to the inclined plane, we find $\dot{\omega} = \Omega \times \omega$ and so $|\dot{\omega}| = \Omega \omega$.

It is easy to see that the velocities of the points lying on the axis of the body can be expressed as peripheral speeds due to Ω and ω as well. In the special case of the CM we'll find them as $v_{\text{CM}} = D \sin \gamma \cdot \omega = D \cos \gamma \cdot \Omega$, thus $\Omega = \omega \tan \gamma$. Consequently, the requested "centripetal" angular acceleration is (we have used that $\tan \gamma = \frac{R_2 - R_1}{2\sqrt{R_1 R_2}}$)

$$|\dot{\omega}| = \omega^2 \tan \gamma = \frac{10}{7} g \frac{R_1^4 + R_2^4}{R_1^5 + R_2^5} \sin \alpha = 180.3 \text{ s}^{-2}.$$

3 We have the equation $\dot{\theta} = \Omega = \omega \tan \gamma$; substituting this into the conservation of energy yields that

$$\frac{8\pi}{3} \rho g \frac{R_1^4 + R_2^4}{R_1 - R_2} \sqrt{R_1 R_2} \cos \theta \sin \alpha + \frac{56\pi}{15} \rho \frac{(R_1^5 + R_2^5) R_1 R_2}{(R_1 - R_2)^2} (\dot{\theta})^2$$

is constant, therefore its time derivative is 0. As the only parameter varying by time is θ , we'll have a differential equation in θ . This equation may be simplified and as θ is very small, we may use first-order approximations. Using these simplifications we'll get

$$5g(R_1^4 + R_2^4) \sin \alpha \cdot \theta = 14 \frac{R_1^5 + R_2^5}{R_1 - R_2} \sqrt{R_1 R_2} \cdot \ddot{\theta};$$

from this harmonic equation we can easily get the angular frequency:

$$\Omega_0 = \sqrt{\frac{5g}{14} \frac{(R_1^4 + R_2^4)(R_2 - R_1)}{(R_1^5 + R_2^5) \sqrt{R_1 R_2}}} \sin \alpha = 5.646 \text{ s}^{-1}.$$

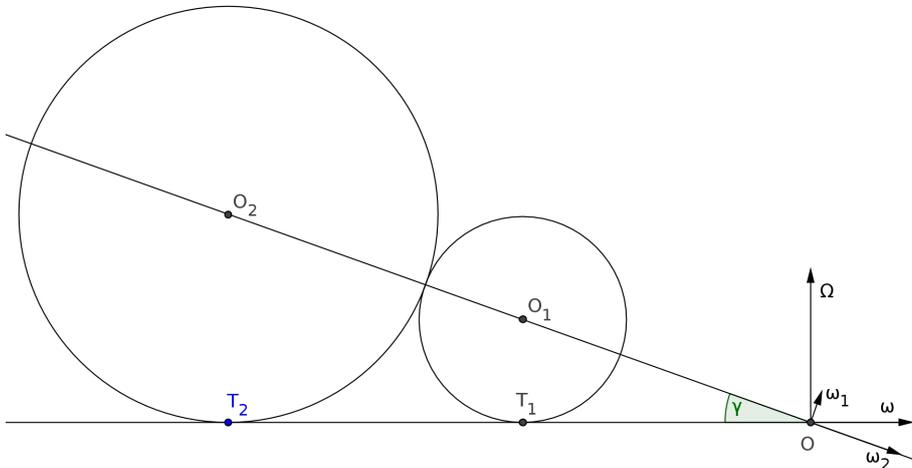


Figure 3. The setup of the system in Part C.